# Reactions of Dimethyl Ether with Atomic Oxygen: A Matrix Isolation and a Quantum Chemical Study 

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#### Abstract

The reaction of dimethyl ether (1) with atomic oxygen generated by photolysis of ozone or $\mathrm{N}_{2} \mathrm{O}$ was examined in low-temperature matrices. The major reaction products are two conformers of methoxymethanol (5). IR absorptions of the products were assigned by isotopic labeling ( ${ }^{18} \mathrm{O}$ and D ) and DFT calculations at the B3LYP/ $6-311++\mathrm{G}(\mathrm{d}, \mathrm{p})$ level of theory. The mechanism of the formation of $\mathbf{5}$, in particular H abstraction from $\mathbf{1}$ by atomic oxygen ( $\mathrm{O}^{3} \mathrm{P}$ and $\mathrm{O}^{1} \mathrm{D}$ ), was investigated using UMP, UCCSD(T), and UDFT. In both the H abstraction and the $\mathrm{O}\left({ }^{( } \mathrm{D}\right)$ insertion reaction, the out-of-plane $\mathrm{C}-\mathrm{H}$ bonds of $\mathbf{1}$ are preferentially attacked since the inplane $\mathrm{C}-\mathrm{H}$ bonds are about $10 \mathrm{kcal} / \mathrm{mol}$ stronger. In the case of a reaction with $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$, an Arrhenius activation energy of $3.5 \mathrm{kcal} / \mathrm{mol}$ is calculated at 298 K , which compares well with an experimental value of 2.85 $\mathrm{kcal} / \mathrm{mol}$. In the exit channel of the reaction, a radical-radical complex between $\mathrm{CH}_{3} \mathrm{CH}_{2} \bullet$ and $\bullet \mathrm{OH}(-2.7$ $\mathrm{kcal} / \mathrm{mol}$ relative to separated products) is found. The latter is the starting point for the formation of $\mathbf{5}$ and helps to rationalize the stereoselectivity of the reaction leading to particular conformations of 5 .


## Introduction

Dimethyl ether (1) and its derivatives have been recently proposed as a diesel fuel substitute. ${ }^{1,2}$ Several properties such as high cetane number, reduction of CO and $\mathrm{NO}_{x}$ emissions, and low-cost one-step synthesis from $\mathrm{C}_{1}$ feedstocks make 1 an attractive diesel fuel. Since the emission of ether $\mathbf{1}$ to the atmosphere might be harmful to the environment, the atmospheric chemistry of $\mathbf{1}$ has been recently studied by several authors. ${ }^{3-6}$ The oxidation of $\mathbf{1}$ in the troposphere is mainly initiated by the reaction with hydroxyl radicals leading to the methoxymethyl radical.

$$
\mathrm{CH}_{3} \mathrm{OCH}_{3}+\mathrm{OH} \rightarrow \mathrm{CH}_{3} \mathrm{OCH}_{2}+\mathrm{H}_{2} \mathrm{O}
$$

The rate of this reaction dictates the atmospheric lifetime of dimethyl ether 1 and influences its greenhouse warming and ozone depletion potential. Once produced, the $\mathrm{CH}_{3} \mathrm{OCH}_{2}$ radical can participate in further degradation in the atmosphere. The reaction of the $\mathrm{CH}_{3} \mathrm{OCH}_{2}$ radical with $\mathrm{O}_{2}$ produces organic peroxy radicals $\mathrm{CH}_{3} \mathrm{OCH}_{2} \mathrm{OO}$, which eventually decompose to formaldehyde. ${ }^{6-7}$

$$
\begin{gathered}
\mathrm{CH}_{3} \mathrm{OCH}_{2}+\mathrm{O}_{2} \rightarrow \mathrm{CH}_{3} \mathrm{OCH}_{2} \mathrm{OO}^{\#} \\
\mathrm{CH}_{3} \mathrm{OCH}_{2} \mathrm{OO}^{\#} \rightarrow \mathrm{CH}_{2} \mathrm{OCH}_{2} \mathrm{OOH}^{\#} \rightarrow 2 \mathrm{HCHO}+\mathrm{OH}
\end{gathered}
$$

Moreover, the radicals are able to react with atmospheric $\mathrm{NO}_{x}$. The above reaction is the key process to understand the generation of secondary pollutants, such as ozone, various carbonyl compounds and organic nitrogen compounds. ${ }^{5}$

In the laboratory, the $\mathrm{CH}_{3} \mathrm{OCH}_{2}$ radical can be produced by hydrogen abstraction from dimethyl ether $\mathbf{1}$ with atomic oxygen. The kinetics of this reaction was investigated by H. LeFevre et

[^0]al. ${ }^{8}$ Alternatively, oxygen atoms might add to $\mathbf{1}$ and thus provide a simple entry to the $\mathrm{C}_{2} \mathrm{H}_{6} \mathrm{O}_{2}$ potential energy surface. Schwarz and co-workers calculated several stationary points of this potential energy surface at the B3LYP/6-311++G (d,p) level of theory. ${ }^{9}$ According to these calculations, the elusive dimethyl ether O-oxide 2 should be $51 \mathrm{kcal} / \mathrm{mol}$ higher in energy than dimethylperoxide 4 , however, separated from 4 by a large activation barrier of $34.9 \mathrm{kcal} / \mathrm{mol}$. The exothermic fragmentation of $\mathbf{2}$ to $\mathrm{CH}_{3} \mathrm{OCH}_{2}$ and OH radicals is predicted to take place with a barrier of only $25.9 \mathrm{kcal} / \mathrm{mol}$, considerably less than the rearrangement to peroxide $\mathbf{4}$, and thus the radical path is the expected major decomposition route. According to the DFT calculations, the reaction of singlet oxygen atoms $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ and $\mathbf{1}$ to yield ether oxide 2 is exothermic by $44.5 \mathrm{kcal} / \mathrm{mol}$, while triplet oxygen atoms $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ are predicted to produce a weakly bound ( 2.8 kcal ) charge transfer complex 3. ${ }^{9}$


In a similar study, Schriver-Mazzuoli et al. ${ }^{10}$ reported on the reaction of $\mathrm{CH}_{4}$ and $\mathrm{CH}_{3} \mathrm{OH}$ with oxygen atoms. These authors found that triplet oxygen atoms $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ are unreactive, while singlet oxygen atoms $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ directly inserted into CH bonds to give $\mathrm{CH}_{3} \mathrm{OH}$ and $\mathrm{HOCH}_{2} \mathrm{OH}$, respectively. A minor route in the oxidation of $\mathrm{CH}_{3} \mathrm{OH}$ was the formation of $\mathrm{CH}_{2} \mathrm{O}$ and $\mathrm{H}_{2} \mathrm{O}$. Peroxides or other products with $\mathrm{O}-\mathrm{O}$ bonds were not observed in this study.

In this work, we investigated the reaction of several isotopomers of dimethyl ether 1 with oxygen atoms under the
conditions of matrix isolation to elucidate the reaction mechanism and to distinguish between the alternatives (i) O -insertion, (ii) O -addition, and (iii) H -abstraction/ OH -addition. Our experimental investigation was supported by appropriate ab initio and density functional theory (DFT) calculations.

## Results and Discussion

Matrix Isolation Studies. Dimethyl ether 1 and two of its isotopomers $\left({ }^{18} \mathrm{O}-\mathbf{1}\right.$ and $\left.\mathrm{d}_{6} \mathbf{- 1}\right)$ were isolated in an argon matrix at 10 K , and the IR spectra were assigned by comparison with calculations at the B3LYP/6-311++G(d,p) level of theory (Table 1). This assignment is in reasonable agreement with an earlier study on crystalline films of $\mathbf{1 . 1}$ The most characteristic absorptions of $\mathbf{1}$ are the asymmetrical COC stretching vibration ( $v_{\text {as }}$ COC, \#8) at $1173 \mathrm{~cm}^{-1}$ and the symmetrical COC stretching vibration ( $v_{\mathrm{s}} \mathrm{COC}, \# 4$ ) at $924 \mathrm{~cm}^{-1}$. As expected, $v_{\mathrm{as}} \mathrm{COC}$ exhibits a large ${ }^{18} \mathrm{O}$ isotopic shift of 0.981 and a comparatively small deuterium isotopic shift of 0.982 . For $v_{\mathrm{s}} \mathrm{COC}$, the ${ }^{18} \mathrm{O}$ isotopic shift is smaller ( 0.987 ), while now a substantial deuterium isotopic shift is observed ( 0.890 ). These spectroscopic data are nicely reproduced by the DFT calculations (Table 1).
$\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ is readily prepared by $\lambda=266 \mathrm{~nm}$ photolysis of matrixisolated $\mathrm{O}_{3}$ (more than $90 \%$ yield of $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ ), ${ }^{12 \mathrm{a}}$ or by $\lambda<200$ nm photolysis of $\mathrm{N}_{2} \mathrm{O} .{ }^{13,14}$ In solid argon the lifetime of $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ is long enough to escape from the primary matrix cage. Intersystem crossing (ISC) finally produces the ground-state $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ which is trapped in interstitial matrix sites. ${ }^{15}$ Triplet oxygen atoms can be directly produced by irradiation of $\mathrm{O}_{3}$ in the near UV or visible region ( $\lambda>360 \mathrm{~nm}$ ) of the spectrum. ${ }^{12 b}$

Dimethyl ether $\mathbf{1}$ was matrix-isolated in argon doped with $0.5-5 \% \mathrm{O}_{3}$ at 10 K and irradiated with $\lambda>420 \mathrm{~nm}(\mathrm{Hg}$ arc lamp with cutoff filter) or $\lambda=248 \mathrm{~nm}$ (KrF Excimer Laser). At high concentrations of $\mathrm{O}_{3}(>1 \%)$, the irradiation resulted in a significant broadening of all IR absorptions, presumably due to the formation of complexes of $\mathbf{1}$ with $\mathrm{O}_{3} \cdot{ }^{16,17}$ With $0.5 \% \mathrm{O}_{3}$, narrow line widths were observed, and therefore this concentration was used throughout the experiments described below. However, since under these conditions only a small fraction of $\mathbf{1}$ is trapped in matrix sites in proximity to $\mathrm{O}_{3}$ molecules and since the diffusion of trapped species at 8 K is slow, the yield of photooxidation products is low ( $<5 \%$ ). On the other hand, the narrow IR line widths allow the spectra to be assigned accurately and experimental and calculated data to be compared.

Irradiation $(\lambda=248 \mathrm{~nm})$ of the matrix until complete disappearance of $\mathrm{O}_{3}$ resulted in the formation of several new IR absorptions (Figure 1), which were stable toward prolonged irradiation. The same set of absorptions was observed when $\mathrm{N}_{2} \mathrm{O}$ was used as a source of oxygen atoms, which indicates that $\mathrm{O}_{3}$ and $\mathrm{O}_{2}$ are not directly involved in this oxidation process. Since the photolysis of $\mathrm{N}_{2} \mathrm{O}$ requires short-wavelength UV irradiation ( $\lambda=193 \mathrm{~nm}$ ) and the resulting IR spectra were less well resolved, $\mathrm{O}_{3}$ was used as the superior source of oxygen atoms in most of the experiments.

By comparison of the new absorptions with the vibrational spectrum of dimethylperoxide 4 reported by Christe, ${ }^{18}$ this compound was definitely excluded as a product of the photooxidation. Traces of methanol and formaldehyde were identified by comparison with authentic matrix-isolated material, while the major IR absorptions could not be assigned to known compounds. In the UV/vis spectrum, a new, broad absorption with a maximum at 300 nm was assigned to the oxidation products. A similar absorption was reported by S. Langer et al. for the UV spectrum of the $\mathrm{CH}_{3} \mathrm{OCH}_{2}$ radical. ${ }^{5}$ However, the corresponding IR absorptions assigned to this radical were not observed in our experiments.
TABLE 1: IR Absorptions of Dimethyl Ether 1 Matrix Isolated in Argon at 10 K and B3LYP/6-311++G(d,p) Vibrational Frequencies of 1, Unscaled

| no. |  | B3LYP/6-311+G(d,p) |  |  |  |  | argon, 10 K |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | sym | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{3}} v,$ | $\begin{gathered} \mathrm{CH}_{3}-18 \mathrm{O}-\mathrm{CH}_{3} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | [ $\nu_{i} / v$ ] | $\begin{aligned} & \mathrm{CD}_{3}-\mathrm{O}-\mathrm{CD}_{3} v, \\ & {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{aligned}$ | [ $v_{i} / \nu$ ] | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{3}} v,$ | $\begin{gathered} \mathrm{CH}_{3}-18 \mathrm{O}-\mathrm{CH}_{3} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | $\left[v_{i} / v\right]$ | $\begin{aligned} & \mathrm{CD}_{3}-\mathrm{O}-\mathrm{CD}_{3} v, \\ & {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{aligned}$ | [ $v_{i} / \nu$ ] | assgnt |
| 1 | $\mathrm{a}_{2}$ | 203 (0) | 203 (0) | 1.0 | 145 (0) | 0.714 |  |  |  |  |  | t $\mathrm{CH}_{3}$ |
| 2 | $\mathrm{b}_{1}$ | 237 (4) | 234 (4) | 0.987 | 184 (4) | 0.776 |  |  |  |  |  | $\mathrm{t} \mathrm{CH}_{3}$ |
| 3 | $\mathrm{a}_{1}$ | 408 (1) | 402 (1) | 0.985 | 344 (1) | 0.843 |  |  |  |  |  | $\delta$ COC |
| 4 | $\mathrm{a}_{1}$ | 936 (29) | 921 (28) | 0.984 | 831 (15) | 0.889 | 924 (35) | 912 (55) | 0.987 | 823 (5) | 0.890 | $\nu_{\mathrm{s}} \mathrm{COC}+\mathrm{rCH}_{3}$ |
| 5 | $\mathrm{b}_{2}$ | 1114 (31) | 1100 (51) | 0.987 | 865 (0) | 0.776 | 1096 (45) | 1084 (97) | 0.989 |  |  | $v \mathrm{COC}+\mathrm{rCH}_{3}$ |
| 6 | $\mathrm{a}_{2}$ | 1155 (0) | 1155 (0) | 1.0 | 877 (0) | 0.759 |  |  |  |  |  |  |
| 7 | $\mathrm{b}_{1}$ | 1186 (4) | 1180 (4) | 0.995 | 937 (6) | 0.794 |  |  |  | 928 (4) |  | r $\mathrm{CH}_{3}$ |
| 8 | $\mathrm{b}_{2}$ | 1192 (69) | 1169 (42) | 0.981 | 1170 (100) | 0.981 | 1173 (100) | 1151 (100) | 0.981 | 1153 (100) | 0.982 | $v \mathrm{COC}+\mathrm{r} \mathrm{CH}_{3}$ |
| 9 | $\mathrm{a}_{1}$ | 1261 (3) | 1251 (2) | 0.992 | 1086 (0) | 0.861 | 1241 (6) | 1236 (6) | 0.996 |  |  | $\delta_{\text {s }} \mathrm{CH}_{3}$ |
| 10 | $\mathrm{b}_{2}$ | 1458 (1) | 1456 (2) | 0.999 | 1077 (1) | 0.739 | 1425 (1) | 1425 (5) | 1.0 |  |  | $\delta_{\text {s }} \mathrm{CH}_{3}$ |
| 11 | $\mathrm{a}_{2}$ | 1478 (0) | 1478 (0) | 1.0 | 1068 (0) | 0.723 |  |  |  |  |  |  |
| 12 | $\mathrm{a}_{1}$ | 1484 (0) | 1484 (0) | 1.0 | 1059 (10) | 0.714 |  |  |  | 1051 (5) |  | $\delta_{\text {as }} \mathrm{CH}_{3}$ |
| 13 | $\mathrm{b}_{1}$ | 1487 (10) | 1487 (10) | 1.0 | 1075 (3) | 0.723 | 1456 (29) | 1456 (34) | 1.0 | 1061 (3) | 0.729 | $\delta_{\text {as }} \mathrm{CH}_{3}$ |
| 14 | $\mathrm{b}_{2}$ | 1494 (10) | 1494 (9) | 1.0 | 1079 (2) | 0.722 | 1461 (15) | 1461 (39) | 1.0 | 1059 (4) | 0.724 | $\delta_{\text {as }} \mathrm{CH}_{3}$ |
| 15 | $\mathrm{a}_{1}$ | 1513 (1) | 1512 (1) | 0.999 | 1163 (4) | 0.769 | 1478 (16) | 1478 (16) | 1.0 |  |  | $\delta_{\text {as }} \mathrm{CH}_{3}$ |
| 16 | $\mathrm{b}_{2}$ | 2958 (41) | 2958 (41) | 1.0 | 2121 (33) | 0.717 | 2820 (50) | 2820 (98) | 1.0 | 2055 (14) | 0.728 | $v \mathrm{~s}^{\mathrm{CH}_{2}}+v \mathrm{CH}$ |
| 17 | $\mathrm{a}_{1}$ | 2970 (52) | 2970 (52) | 1.0 | 2133 (27) | 0.718 | 2869 (15) | 2869 (27) | 1.0 |  |  | $v \mathrm{SH}_{2}+v \mathrm{CH}$ |
| 18 | $\mathrm{b}_{1}$ | 3003 (100) | 3003 (100) | 1.0 | 2226 (59) | 0.741 | 2891 (43) | 2891 (95) | 1.0 | 2190 (16) | 0.757 | $v_{\text {as }} \mathrm{CH}_{2}$ |
| 19 | $\mathrm{a}_{2}$ | 3008 (0) | 3008 (0) | 1.0 | 2232 (0) | 0.742 | 2934 (13) | 2934 (48) | 1.0 |  |  | $1478+1456$ |
| 20 | $\mathrm{b}_{2}$ | 3113 (22) | 3113 (22) | 1.0 | 2303 (8) | 0.740 | 2939 (19) | 2939 (43) | 1.0 | 2250 (5) | 0.765 | $v \mathrm{CH}+\nu_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 21 | $\mathrm{a}_{1}$ | 3114 (14) | 3114 (14) | 1.0 | 2307 (13) | 0.741 | 2986 (24) | 2986 (54) | 1.0 | 2310 (1) | 0.774 | $v \mathrm{CH}+\nu_{\mathrm{s}} \mathrm{CH}_{2}$ |



Figure 1. (a) IR difference spectrum showing the oxidation of $\mathbf{1}$ with atomic O in argon at 10 K after irradiation of the matrix of $\mathbf{1}$ and ozone ( $0.8 \%$ ) with $\lambda=248 \mathrm{~nm}$; bottom part, bands disappearing; bands of $\mathbf{5 a}$ and $\mathbf{5 b}$ appearing on irradiation. (b) Calculated spectrum of $\mathbf{5 a}$ at the B3LYP/6311+ $+G(d, p)$ level of theory, unscaled. (c) Calculated spectrum of $\mathbf{5 b}$ at the B3LYP/6311++G(d,p) level of theory, unscaled.

The IR spectrum of the products obtained by 248 nm irradiation of the $1 / \mathrm{O}_{3}$ mixture clearly shows OH stretching vibrations at 3631 and $3641 \mathrm{~cm}^{-1}$, which on deuteration are red-shifted by more than $900 \mathrm{~cm}^{-1}$ (Tables 2 and 3). By comparing the IR spectrum with data calculated at the B3LYP/ $6-311++G(d, p)$ level of theory, two conformers of methoxymethanol ${ }^{5}$ were identified as the main products of the oxidation (Tables 4 and 5). The hemi-aceta ${ }^{5}$ is the main product of the equilibrium reaction between formaldehyde and methanol. ${ }^{19}$ In substance, $\mathbf{5}$ is not stable and IR spectra were unknown, so far, while ${ }^{1} \mathrm{H}$ NMR spectra of $\mathbf{5}$ could be recorded at low temperature. Since under the conditions of matrix isolation the cleavage of $\mathbf{5}$ is excluded, alternative mechanisms have to be considered to rationalize the formation of traces of methanol and formaldehyde in our experiments. One route could be the photochemical cleavage of the CO bond of 5 into the $\mathrm{CH}_{3} \mathrm{O}$ and $\mathrm{CH}_{2} \mathrm{OH}$ radicals. The $\mathrm{CH}_{3} \mathrm{O}$ radical rearranges rapidly into the more stable $\mathrm{CH}_{2} \mathrm{OH}$ radical, ${ }^{20-23}$ which subsequently disproportionates into $\mathrm{CH}_{2} \mathrm{O}$ and $\mathrm{CH}_{3} \mathrm{OH} .{ }^{24}$

By oxidizing $\mathbf{1},{ }^{18} \mathrm{O}-\mathbf{1}$, and $\mathrm{d}_{6}-\mathbf{1}$ with both ${ }^{16} \mathrm{O}_{3}$ and ${ }^{18} \mathrm{O}_{3}$, six isotopomers of 5 were obtained, which allowed a reliable assignment of the IR absorptions. The most characteristic IR absorptions of 5a are the weak absorption of the combination vibration $(v \mathrm{COC}+\delta \mathrm{OH}, \# 5)$ at $576 \mathrm{~cm}^{-1}$, the COC stretching vibration ( $v_{\mathrm{s}}$ COC, \#6) at $936 \mathrm{~cm}^{-1}$, the strong absorption at $1020 \mathrm{~cm}^{-1}(v \mathrm{COH}, \# 7)$, a combination vibration $(v \mathrm{COC}+$ $\mathrm{CH}_{2}$ scis, \#8) at $1044 \mathrm{~cm}^{-1}$, and the strongest absorption ( $v$ $\mathrm{COC}+\omega \mathrm{CH} 2, \# 9)$ at $1125 \mathrm{~cm}^{-1}$. Vibration \#5 shows a large ${ }^{18} \mathrm{O}$ isotopic shift and a relatively small deuterium isotopic shift (0.925). The symmetrical COC stretching vibration (\#6) is blueshifted compared to the corresponding vibration in $\mathbf{1}$ and shows similar isotopic ${ }^{18} \mathrm{O}$ or deuterium isotopic shifts. Vibration \#7 exhibits both large ${ }^{18} \mathrm{O}(0.982)$ and deuterium ( 0.959 ) isotopic shifts on isotopic substitution of the OH group. Vibration \#8 is located in the same area as a strong absorption of $\mathrm{O}_{3}$ and thus not observed in the ${ }^{16} \mathrm{O}$ experiments, but by reaction with ${ }^{18} \mathrm{O}_{3}$ this absorption is clearly observed. \#9 shows a large ${ }^{18} \mathrm{O}$ isotopic shift for the isotopic labeling of oxygen by ether $\mathbf{1}(0.981)$. The
$\mathrm{CH}_{2}$ wag vibration (\#14) shows a very large deuterium isotopic shift ( 0.788 ) and is the most intense absorption in the perdeuterated 5a.

The second conformer $\mathbf{5 b}$ is calculated to be higher in energy by $2.1 \mathrm{kcal} / \mathrm{mol}$. The characteristic absorptions are the symmetrical COC stretching vibration ( $v_{\mathrm{s}} \mathrm{COC}, \# 6$ ) at $940 \mathrm{~cm}^{-1}$, the very strong combination vibration at $1019 \mathrm{~cm}^{-1}(\nu \mathrm{COH}$, \#7), and the combination vibration at $1119 \mathrm{~cm}^{-1}(v \mathrm{COC}+\omega$ CH 2 , \#9, Table 3). Isotopic labeling at the OH group results in the expected large ${ }^{18} \mathrm{O}(0.985)$ and deuterium ( 0.888 ) isotopic shift of vibration \#6. For the strongest absorption, \#7, labeling of the OH group leads to a small ${ }^{18} \mathrm{O}(0.980)$ and a relatively large deuterium isotopic shift ( 0.964 ). The combination vibration \#9 shows a very large ${ }^{18} \mathrm{O}$ isotopic shift on labeling of the ether oxygen atom (0.974). As in $\mathbf{5 a}$, the most intense absorption after complete deuterium labeling of $\mathbf{5 b}$ with a very large isotopic shift ( 0.784 ) is assigned to the $\mathrm{CH}_{2}$ wag vibration (\#14). The position of the OH stretching vibrations (\#24) is similar for both conformers of 5 (3631 and $3641 \mathrm{~cm}^{-1}$ for $\mathbf{5 a}$ and $\mathbf{5 b}$, respectively). Both vibrations show a small (0.997) ${ }^{18} \mathrm{O}$ and a very large (0.738) deuterium isotopic shift.

Conformational Behavior of $\mathbf{C H}_{\mathbf{3}} \mathbf{O C H}_{\mathbf{2}} \mathbf{O H}$ (5). There are five forms of $\mathbf{5}$ in which all vicinal bonds at the CO bonds are staggered: 5-1 (ap,ap), 5-2 (ap, +sc$), \mathbf{5 - 3}(+\mathrm{sc}, \mathrm{ap}), \mathbf{5 - 4}(+\mathrm{sc},+\mathrm{sc})$, $\mathbf{5 - 5}(+\mathrm{sc},-\mathrm{sc})$ (Scheme 1). Geometry optimizations at the B3LYP/6-311++G(d,p) level of theory show that only three of these forms are close to stable conformers (5-4 $\boldsymbol{\mathbf { 5 a } , \mathbf { 5 - 5 } \rightarrow}$ $\mathbf{5 b}, \mathbf{5 - 2} \rightarrow \mathbf{5 e}$ ) while forms 5-1 (ap,ap) and 5-3 ( $+\mathrm{sc}, \mathrm{ap}$ ) are equal or close to transition states (TS) on the conformational energy surface (CSE) spanned by the two rotational angles $\tau_{1}$ $=\angle \mathrm{MeO}-\mathrm{CO}$ and $\tau_{2}=\angle \mathrm{HO}-\mathrm{CO}$ of 5 (Scheme 1). Conformation 5a occupies the global minimum of the CES at $\tau_{1}=$ $68.9^{\circ}$ and $\tau_{2}=64.9^{\circ}$ while $\mathbf{5 b}\left(2.1 \mathrm{kcal} / \mathrm{mol}\right.$ relative to $\mathbf{5 a}$; $\tau_{1}$ $=69.9^{\circ}, \tau_{2}=-85.1^{\circ}$, Scheme 1a) and 5e ( $2.2 \mathrm{kcal} / \mathrm{mol} ; \tau_{1}=$ $180^{\circ}, \tau_{2}=56.4^{\circ}$, Scheme 1a) correspond to local minima. In a previous investigation of the CES of 5, Jeffrey, Pople, and Radom found conformational minima similar to 5a and 5e but missed 5b since their work was based on RHF/4-31G theory

TABLE 2: IR Absorptions of Methoxymethanol (5a) Isolated in Argon at 10 K

| No. | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{2} \mathrm{OH}} v,$ | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{2} \mathrm{O}^{18} \mathrm{H} v,}$ | [ $v_{i} / v$ ] | $\underset{\left[\mathrm{cm}^{-1}\right]\left(\mathrm{I}_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}^{18}-\mathrm{CH}_{2} \mathrm{OH}} v,$ | [ $v_{i} / v$ ] | $\underset{\left[\mathrm{cm}^{-1}\right]\left(\mathrm{I}_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}^{18}-\mathrm{CH}_{2} \mathrm{O}^{18} \mathrm{H} v,}$ | [ $v_{i} / v$ ] | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CD}_{3}-\mathrm{O}-\mathrm{CD}_{2} \mathrm{OD}} v,$ | [ $v_{i} / v$ ] | $\begin{gathered} \mathrm{CD}_{3}-\mathrm{O}-\mathrm{CD}_{2} \mathrm{O}^{18} \mathrm{D} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | [ $v_{i} / v$ ] | assgnt |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 5 | 576 (14) | 566 (7) | 0.983 | 562 (8) | 0.976 | 554 (8) | 0.962 | 533 (12) | 0.925 | 525 (6) | 0.909 | $v \mathrm{COC}+\delta \mathrm{OH}$ |
| 6 | 936 (20) | 936 (31) | 1.0 | 922 (32) | 0.985 | 921 (26) | 0.984 | 833 (28) | 0.900 | 833 (23) | 0.900 | $\nu_{\mathrm{s}} \mathrm{COC}$ |
| 7 | 1020 (94) | 1002 (100) | 0.982 | 1119 (80) | 0.999 | 1001 (100) | 0.981 | 978 (25) | 0.959 | 958 (43) | 0.939 | $v \mathrm{COH}$ |
| 8 | 1044 (-) | 1039 (47) | 0.995 |  |  | 1033 (56) | 0.990 |  |  | 809 (5) | 0.774 | $v \mathrm{COC}+\mathrm{scisCH}_{2}$ |
| 9 | 1125 (100) | 1123 (36) | 0.998 | 1104 (44) | 0.981 | 1103 (24) | 0.986 |  |  |  |  | $v$ COC $+\omega \mathrm{CH}_{2}$ |
| 10 | 1150 (2) |  |  |  |  |  |  |  |  |  |  | $-$ |
| 11 | 1187 (34) | 1185 (24) | 0.999 | 1177 (20) | 0.992 | 1177 (12) | 0.992 |  |  |  |  | $\omega \mathrm{CH}_{3}+\omega \mathrm{CH}_{2}$ |
| 12 | 1286 (9) | 1285 (7) | 0.999 | 1280 (6) | 0.995 | 1279 (5) | 0.995 |  |  | 1027 (8) | 0.798 | $\omega \mathrm{CH}_{3}+\omega \mathrm{CH}_{2}$ |
| 13 | 1355 (10) | 1350 (9) | 0.996 | 1354 (8) | 0.999 | 1348 (8) | 0.995 | 1080 (4) | 0.797 | 1078 (6) | 0.795 | $\delta \mathrm{CH}_{2}$ |
| 14 | 1406 (3) | 1406 (8) | 1.0 | 1402 (11) | 0.997 | 1401 (9) | 0.998 | 1109 (100) | 0.788 | 1108 (100) | 0.787 | wag $\mathrm{CH}_{2}$ |
| 15 | 1444 (4) | 1443 (4) | 0.999 | 1444 (4) | 1.0 | 1444 (4) | 1.0 |  |  |  |  | $\delta \mathrm{CH}_{3}$ |
| 16 | 1452 (7) | 1452 (5) | 1.0 | 1452 (5) | 1.0 | 1451 (5) | 1.0 |  |  |  |  |  |
| 17 | 1470 (6) | 1470 (5) | 1.0 | 1470 (5) | 1.0 | 1470 (6) | 1.0 |  |  |  |  |  |
| 18 |  |  |  |  |  |  |  | 1142 (18) |  | 1131 (18) |  | $\delta \mathrm{CH}_{3}$ |
| 19 | 2878 (13) | 2878 (13) | 1.0 | 2878 (13) | 1.0 | 2878 (13) | 1.0 | 2062 (22) | 0.716 | 2062 (22) | 0.716 | $\nu_{\text {as }} \mathrm{CH}_{2}$ |
| 20 | 2929 (15) | 2929 (15) | 1.0 | 2929 (15) | 1.0 | 2929 (15) | 1.0 |  |  |  |  | $v_{\text {as }} \mathrm{CH}_{2}$ |
| 21 | 2964 (13) | 2964 (13) | 1.0 | 2964 (13) | 1.0 | 2964 (13) | 1.0 | 2260 (19) | 0.762 | 2260 (19) | 0.762 | $v \mathrm{CH}+v_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 22 |  |  |  |  |  |  |  | 2346 (16) |  | 2346 (16) |  | $v \mathrm{CH}+v_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 23 |  |  |  |  |  |  |  |  |  |  |  |  |
| 24 | 3631 (18) | 3622 (16) | 0.997 | 3631 (19) | 1.0 | 3622 (16) | 0.997 | 2682 (10) | 0.737 | 2664 (10) | 0.733 | $v \mathrm{O}-\mathrm{H}$ |

${ }^{a}$ Relative intensity based on the strongest absorption.

TABLE 3: IR-Absorptions of Methoxymethanol (5b) Isolated in Argon at 10 K

| No. | $\begin{gathered} \mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{2} \mathrm{OH} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{2} \mathrm{O}^{18} \mathrm{H} v,}$ | [ $v_{i} / \nu$ ] | $\begin{gathered} \mathrm{CH}_{3}-\mathrm{O}^{18}-\mathrm{CH}_{2} \mathrm{OH} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | $\left[\nu_{i} / \nu\right]$ | $\begin{gathered} \mathrm{CH}_{3}-\mathrm{O}^{18}-\mathrm{CH}_{2} \mathrm{O}^{18} \mathrm{H} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | $\left[v_{i} / v\right]$ | $\begin{gathered} \mathrm{CD}_{3}-\mathrm{O}-\mathrm{CD}_{2} \mathrm{OD} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | $\left[v_{i} / v\right]$ | $\begin{gathered} \mathrm{CD}_{3}-\mathrm{O}-\mathrm{CD}_{2} \mathrm{O}^{18} \mathrm{D} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | $\left[v_{i} / v\right]$ | assgnt |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 5 | 571 (6) | 564 (8) | 0.988 | 558 (8) | 0.977 | 550 (4) | 0.963 |  |  |  |  | $v \mathrm{COC}+\delta \mathrm{OH}$ |
| 6 | 940 (37) | 940 (43) | 1.0 | 926 (38) | 0.985 | 926 (26) | 0.985 | 835 (31) | 0.888 | 835 (45) | 0.888 | $\nu_{\mathrm{s}} \mathrm{COC}$ |
| 7 | 1019 (100) | 999 (100) | 0.980 | 1018 (100) | 0.998 | 998 (100) | 0.979 | 982 (25) | 0.964 | 959 (33) | 0.941 | $v \mathrm{COH}$ |
| 8 | 1064 (17) |  |  |  |  |  |  |  |  |  |  | $v \mathrm{COC}+\mathrm{scisCH}_{2}$ |
| 9 | 1119 (55) | 1119 (53) | 1.0 | 1090 (86) | 0.974 | 1090 (79) | 0.974 |  |  |  |  | $v \mathrm{COC}+\omega \mathrm{CH}_{2}$ |
| 10 | 1143 (2) |  |  |  |  |  |  |  |  |  |  |  |
| 11 | 1186 (43) | 1185 (24) | 0.999 | 1177 (20) | 0.992 | 1177 (13) | 0.992 |  |  |  |  | $\omega \mathrm{CH}_{3}+\omega \mathrm{CH}_{2}$ |
| 12 | 1285 (12) | 1280 (6) | 0.996 | 1280 (7) | 0.996 | 1277 (4) | 0.994 |  |  |  |  | $\omega \mathrm{CH}_{3}+\omega \mathrm{CH}_{2}$ |
| 13 | 1356 (15) | 1350 (7) | 0.996 | 1355 (8) | 0.999 | 1346 (5) | 0.993 |  |  |  |  | $\delta \mathrm{CH}_{2}$ |
| 14 | 1414 (5) | 1414 (8) | 1.0 | 1412 (2) | 0.998 | 1411 (3) | 0.997 | 1109 (100) | 0.784 | 1109 (100) | 0.784 | wag CH2 |
| 15 | 1444 (6) | 1444 (4) | 1.0 | 1444 (6) | 1.0 | 1444 (4) | 1.0 |  |  |  |  | $\delta \mathrm{CH}_{3}$ |
| 16 | 1452 (6) | 1452 (5) | 1.0 | 1452 (6) | 1.0 | 1452 (5) | 1.0 |  |  |  |  | $\delta \mathrm{CH}_{3}$ |
| 17 | 1470 (10) | 1470 (5) | 1.0 | 1470 (10) | 1.0 | 1470 (5) | 1.0 |  |  |  |  | $\delta \mathrm{CH}_{3}$ |
| 18 |  |  |  |  |  |  |  | 1131 (25) |  | 1131 (25) |  | $\delta \mathrm{CH}_{3}$ |
| 19 | 2824 (14) | 2824 (14) | 1.0 | 2824 (14) | 1.0 | 2824 (14) | 1.0 | 2068 (22) | 0.732 | 2068 (22) | 0.732 | $v_{\text {as }} \mathrm{CH}_{2}$ |
| 20 |  |  |  |  |  |  |  | 2189 (34) |  | 2189 (34) |  | $\nu_{\text {as }} \mathrm{CH}_{2}$ |
| 21 | 2946 (12) | 2946 (12) | 1.0 | 2946 (12) | 1.0 | 2946 (12) | 1.0 | 2305 (11) | 0.782 | 2305 (11) | 0.782 | $v \mathrm{CH}+\nu_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 22 | 3006 (14) | 3006 (14) | 1.0 | 3006 (14) | 1.0 | 3006 (14) | 1.0 |  |  |  |  | $v \mathrm{CH}+v_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 23 |  |  |  |  |  |  |  |  |  |  |  |  |
| 24 | 3641 (11) | 3629 (7) | 0.997 | 3641 (11) | 1.0 | 3629 (7) | 0.997 | 2687 (6) | 0.738 | 2670 (5) | 0.733 | $v \mathrm{O}-\mathrm{H}$ |

TABLE 4: Calculated Vibrational Frequencies of Methoxymethanol (5a) at the B3LYP/6-311++G(d,p) Level of Theory, Unscaled

| No. | $\begin{gathered} \mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{2} \mathrm{OH} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{2} \mathrm{O}^{18} \mathrm{H} v,}$ | [ $v_{i} / v$ ] | $\begin{gathered} \mathrm{CH}_{3}-\mathrm{O}^{18}-\mathrm{CH}_{2} \mathrm{OH} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | [ $v_{i} / \nu$ ] | $\begin{gathered} \mathrm{CH}_{3}-\mathrm{O}^{18}-\mathrm{CH}_{2} \mathrm{O}^{18} \mathrm{H} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | [ $v_{i} / v$ ] | $\begin{gathered} \mathrm{CD}_{3}-\mathrm{O}-\mathrm{CD}_{2} \mathrm{OD} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | [ $v_{i} / v$ ] | $\begin{gathered} \mathrm{CD}_{3}-\mathrm{O}-\mathrm{CD}_{2} \mathrm{O}^{18} \mathrm{D} v, \\ {\left[\mathrm{~cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}} \end{gathered}$ | [ $v_{i} / v$ ] | assgnt |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 144 (2) | 143 (2) | 0.993 | 143 (2) | 0.993 | 142 (2) | 0.986 | 107 (0) | 0.743 | 107 (0) | 0.743 | t CH |
| 2 | 191 (4) | 189 (3) | 0.990 | 190 (5) | 0.995 | 188 (3) | 0.984 | 162 (3) | 0.835 | 160 (3) | 0.838 | $\mathrm{t} \mathrm{CH}_{3}$ |
| 3 | 358 (44) | 356 (35) | 0.994 | 355 (42) | 0.992 | 353 (32) | 0.986 | 275 (35) | 0.768 | 273 (36) | 0.763 | tOH |
| 4 | 394 (36) | 393 (31) | 0.997 | 390 (42) | 0.990 | 389 (33) | 0.987 | 325 (3) | 0.825 | 324 (4) | 0.822 |  |
| 5 | 583 (18) | 573 (14) | 0.983 | 571 (20) | 0.979 | 561 (14) | 0.962 | 536 (11) | 0.919 | 527 (10) | 0.904 | $v \mathrm{COC}+\delta \mathrm{OH}$ |
| 6 | 925 (23) | 925 (18) | 1.0 | 913 (24) | 0.987 | 912 (17) | 0.986 | 834 (13) | 0.902 | 834 (13) | 0.902 | $\nu_{\mathrm{s}} \mathrm{COC}$ |
| 7 | 1024 (100) | 1002 (100) | 0.978 | 1024 (81) | 1.0 | 1002 (100) | 0.978 | 979 (43) | 0.956 | 956 (44) | 0.934 | $v \mathrm{COH}$ |
| 8 | 1037 (56) | 1035 (23) | 0.998 | 1031 (100) | 0.994 | 1029 (34) | 0.992 | 789 (7) | 0.761 | 786 (9) | 0.758 | $v \mathrm{COC}+$ scis $\mathrm{CH}_{2}$ |
| 9 | 1134 (81) | 1133 (70) | 0.999 | 1112 (75) | 0.981 | 1110 (63) | 0.979 | 1069 (1) | 0.945 | 1068 (1) | 0.942 | $\nu \mathrm{COC}+\omega \mathrm{CH}_{2}$ |
| 10 | 1171 (2) | 1171 (2) | 1.0 | 1168 (1) | 0.997 | 1168 (1) | 0.997 | 884 (3) | 0.755 | 883 (4) | 0.754 | $\omega \mathrm{CH}_{2}$ |
| 11 | 1200 (16) | 1200 (14) | 1.0 | 1192 (7) | 0.993 | 1192 (5) | 0.993 | 915 (1) | 0.763 | 915 (1) | 0.763 | $\omega \mathrm{CH}_{3}+\omega \mathrm{CH}_{2}$ |
| 12 | 1299 (9) | 1298 (7) | 0.999 | 1292 (9) | 0.995 | 1291 (7) | 0.994 | 1025 (3) | 0.798 | 1022 (3) | 0.787 | $\omega \mathrm{CH}_{3}+\omega \mathrm{CH}_{2}$ |
| 13 | 1379 (5) | 1373 (4) | 0.996 | 1378 (5) | 0.999 | 1372 (4) | 0.995 | 1094 (10) | 0.793 | 1090 (10) | 0.790 | $\omega \mathrm{CH}_{3}+\delta \mathrm{CH}_{2}$ |
| 14 | 1433 (13) | 1432 (11) | 0.999 | 1431 (15) | 0.999 | 1430 (12) | 0.998 | 1113 (100) | 0.746 | 1112 (100) | 0.745 | wag $\mathrm{CH}_{3}$ |
| 15 | 1473 (6) | 1473 (4) | 1.0 | 1473 (6) | 1.0 | 1473 (4) | 1.0 | 1157 (23) | 0.785 | 1156 (22) | 0.785 | $\delta \mathrm{CH}_{3}$ |
| 16 | 1486 (2) | 1486 (2) | 1.0 | 1485 (3) | 1.0 | 1485 (2) | 1.0 | 1073 (1) | 0.722 | 1073 (1) | 0.722 | $\delta \mathrm{CH}_{3}$ |
| 17 | 1506 (2) | 1506 (1) | 1.0 | 1506 (2) | 1.0 | 1506 (1) | 1.0 | 1086 (3) | 0.721 | 1086 (3) | 0.721 | $\delta \mathrm{CH}_{3}$ |
| 18 | 1511 (3) | 1511 ((2) | 1.0 | 1511 (2) | 1.0 | 1511 (2) | 1.0 | 1153 (35) | 0.763 | 1151 (37) | 0.763 | $\delta \mathrm{CH}_{3}$ |
| 19 | 2994 (31) | 2994 (26) | 1.0 | 2994 (33) | 1.0 | 2994 (26) | 1.0 | 2149 (18) | 0.717 | 2149 (18) | 0.717 | $\nu_{\text {as }} \mathrm{CH}_{2}$ |
| 20 | 3017 (39) | 3017 (32) | 1.0 | 3017 (41) | 1.0 | 3017 (32) | 1.0 | 2189 (23) | 0.725 | 2189 (24) | 0.725 | $v_{\text {as }} \mathrm{CH}_{2}$ |
| 21 | 3064 (20) | 3064 (17) | 1.0 | 3064 (22) | 1.0 | 3064 (17) | 1.0 | 2269 (12) | 0.741 | 2269 (13) | 0.741 | $v \mathrm{CH}+v_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 22 | 3073 (25) | 3073 (21) | 1.0 | 3073 (27) | 1.0 | 3073 (21) | 1.0 | 2289 (13) | 0.745 | 2289 (14) | 0.745 | $v \mathrm{CH}+\nu_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 23 | 3125 (12) | 3125 (10) | 1.0 | 3125 (13) | 1.0 | 3125 (10) | 1.0 | 2317 (6) | 0.741 | 2317 (7) | 0.741 | $v \mathrm{CH}+\nu_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 24 | 3815 (20) | 3802 (16) | 0.997 | 3815 (21) | 1.0 | 3802 (16) | 0.997 | 2777 (10) | 0.728 | 2759 (10) | 0.723 | $v \mathrm{O}-\mathrm{H}$ |

${ }^{a}$ Relative intensity based on the strongest absorption.
TABLE 5: Calculated Vibrational Frequencies of Methoxymethanol (5b) at the B3LYP/6-311++G(d,p) level of Theory, Unscaled

| No. | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{2} \mathrm{OH}} v,$ | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}-\mathrm{CH}_{2} \mathrm{O}^{18} \mathrm{H} v,}$ | [ $v_{i} / v$ ] | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}^{18}-\mathrm{CH}_{2} \mathrm{OH} v,}$ | [ $v_{i} / v$ ] | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CH}_{3}-\mathrm{O}^{18}-\mathrm{CH}_{2} \mathrm{O}^{18} \mathrm{H} v,}$ | [ $v_{i} / \nu$ ] | $\underset{\left[\mathrm{cm}^{-1}\right]\left(I_{\mathrm{rel}}\right)^{a}}{\mathrm{CD}_{3}-\mathrm{O}-\mathrm{CD}_{2} \mathrm{OD}} v,$ | [ $v_{i} / v$ ] | $\underset{\left[\mathrm{cm}^{-1}\right]\left(\mathrm{I}_{\mathrm{rel}}\right)^{a}-\mathrm{O}-\mathrm{CD}_{3} \mathrm{O}^{18} \mathrm{D} v,}{ }$ | [ $v_{i} / v$ ] | assgnt |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 136 (3) | 134 (3) | 0.985 | 135 (3) | 0.992 | 133 (3) | 0.978 | 111 (2) | 0.816 | 111 (2) | 0.816 | t CH |
| 2 | 194 (2) | 193 (2) | 0.995 | 193 (2) | 0.995 | 192 (2) | 0.990 | 151 (5) | 0.778 | 149 (5) | 0.768 | $\mathrm{t} \mathrm{CH}_{3}$ |
| 3 | 303 (32) | 303 (33) | 1.0 | 303 (32) | 1.0 | 302 (32) | 0.997 | 223 (26) | 0.736 | 222 (27) | 0.732 | t OH |
| 4 | 392 (11) | 389 (11) | 0.992 | 387 (12) | 0.987 | 384 (12) | 0.979 | 331 (6) | 0.844 | 329 (6) | 0.839 |  |
| 5 | 569 (3) | 560 (2) | 0.984 | 555 (2) | 0.975 | 546 (2) | 0.959 | 538 (5) | 0.945 | 528 (5) | 0.928 | $v \mathrm{COC}+\delta \mathrm{OH}$ |
| 6 | 929 (15) | 928 (14) | 0.999 | 917 (15) | 0.987 | 917 (14) | 0.987 | 838 (20) | 0.902 | 838 (21) | 0.902 | $\nu_{\mathrm{s}} \mathrm{COC}$ |
| 7 | 1020 (100) | 1001 (100) | 0.981 | 1018 (100) | 0.998 | 999 (100) | 0.979 | 976 (51) | 0.957 | 954 (54) | 0.935 | $\nu \mathrm{COH}$ |
| 8 | 1041 (5) | 1037 (2) | 0.996 | 1040 (6) | 0.999 | 1036 (2) | 0.995 | 770 (9) | 0.740 | 769 (11) | 0.739 | $v \mathrm{COC}+$ scis $\mathrm{CH}_{2}$ |
| 9 | 1122 (44) | 1121 (47) | 0.999 | 1098 (48) | 0.979 | 1096 (51) | 0.977 | 1072 (0) | 0.955 | 1072 (0) | 0.955 | $v \mathrm{COC}+\omega \mathrm{CH}_{2}$ |
| 10 | 1169 (2) | 1168 (2) | 0.999 | 1164 (1) | 0.996 | 1164 (1) | 0.996 | 886 (1) | 0.758 | 886 (1) | 0.758 | $\omega \mathrm{CH}_{2}$ |
| 11 | 1205 (27) | 1205 (28) | 1.0 | 1194 (16) | 0.991 | 1194 (16) | 0.991 | 913 (4) | 0.758 | 912 (5) | 0.757 | $\omega \mathrm{CH}_{3}+\omega \mathrm{CH}_{2}$ |
| 12 | 1304 (9) | 1303 (9) | 0.999 | 1297 (9) | 0.994 | 1296 (9) | 0.994 | 1021 (6) | 0.783 | 1017 (5) | 0.780 | $\omega \mathrm{CH}_{3}+\omega \mathrm{CH}_{2}$ |
| 13 | 1377 (3) | 1371 (2) | 0.996 | 1377 (3) | 1.0 | 1371 (2) | 0.996 | 1094 (6) | 0.794 | 1091 (0) | 0.792 | $\omega \mathrm{CH}_{3}+\delta \mathrm{CH}_{2}$ |
| 14 | 1437 (6) | 1436 (6) | 0.999 | 1434 (7) | 0.998 | 1433 (7) | 0.997 | 1111 (100) | 0.773 | 1111 (100) | 0.773 | wag $\mathrm{CH}_{3}$ |
| 15 | 1471 (2) | 1471 (2) | 1.0 | 1471 (2) | 1.0 | 1471 (2) | 1.0 | 1164 (30) | 0.791 | 1163 (29) | 0.791 | $\delta \mathrm{CH}_{3}$ |
| 16 | 1488 (1) | 1488 (1) | 1.0 | 1488 (1) | 1.0 | 1488 (1) | 1.0 | 1076 (7) | 0.723 | 1075 (10) | 0.723 | $\delta \mathrm{CH}_{3}$ |
| 17 | 1499 (0) | 1499 (0) | 1.0 | 1499 (0) | 1.0 | 1499 (0) | 1.0 | 1088 (12) | 0.726 | 1085 (11) | 0.724 | $\delta \mathrm{CH}_{3}$ |
| 18 | 1511 (3) | 1511 (3) | 1.0 | 1511 (3) | 1.0 | 1511 (3) | 1.0 | 1147 (41) | 0.759 | 1146 (45) | 0.759 | $\delta \mathrm{CH}_{3}$ |
| 19 | 2955 (28) | 2955 (29) | 1.0 | 2955 (28) | 1.0 | 2955 (29) | 1.0 | 2134 (20) | 0.722 | 2134 (21) | 0.722 | $\nu_{\text {as }} \mathrm{CH}_{2}$ |
| 20 | 2973 (27) | 2973 (27) | 1.0 | 2973 (27) | 1.0 | 2973 (27) | 1.0 | 2157 (34) | 0.726 | 2157 (35) | 0.726 | $v_{\text {as }} \mathrm{CH}_{2}$ |
| 21 | 3033 (19) | 3033 (19) | 1.0 | 3033 (19) | 1.0 | 3033 (19) | 1.0 | 2247 (17) | 0.741 | 2247 (18) | 0.741 | $v \mathrm{CH}+v_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 22 | 3121 (8) | 3121 (8) | 1.0 | 3121 (8) | 1.0 | 3121 (8) | 1.0 | 2310 (6) | 0.740 | 2310 (6) | 0.740 | $v \mathrm{CH}+\nu_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 23 | 3128 (8) | 3128 (9) | 1.0 | 3128 (8) | 1.0 | 3128 (8) | 1.0 | 2318 (11) | 0.741 | 2318 (12) | 0.741 | $v \mathrm{CH}+\nu_{\mathrm{s}} \mathrm{CH}_{2}$ |
| 24 | 3842 (11) | 3830 (11) | 0.997 | 3842 (11) | 1.0 | 3830 (11) | 0.997 | 2798 (10) | 0.728 | 2780 (10) | 0.723 | $v \mathrm{O}-\mathrm{H}$ |

## SCHEME $1^{a}$


(a)

-1, (ap,ap)
$5 f$


5-2, (ap,+sc) $5 \mathbf{5}$


5b
(b)

${ }^{a}$ (a) Conformations of $\mathbf{5 - 1}$ to $\mathbf{5 - 5}$ with perfect staggering at the CO bonds with rotational angles $\tau_{1}$ and $\tau_{2}$ indicated. For the ( $\mathrm{sp}, \mathrm{sp}$ ) form, $\tau_{1}=$ $\tau_{2}=0^{\circ}$. Lower numbers correspond to stationary points found on the CES of $\mathbf{5}$. (b) Flip-flop rotation of $\mathbf{5}$ (inner circle) as described by the conformers encountered in this process. Relative energies calculated at the B3LYP/6-311+ $+\mathrm{G}(\mathrm{d}, \mathrm{p})$ level are also given.
and a rigid rotor model defined by standard geometrical parameters. ${ }^{25}$ Such an approach can provide a reasonable account of the topology of the CES; however, the location of stationary points and values of calculated rotational barriers obtained in this way are questionable.

The stability of $\mathbf{5 a}$ and $\mathbf{5 b}$ is caused by anomeric delocalization of one of the electron lone pairs at oxygen into the vicinal CO bond, which will be best accomplished if the OH and OMe bonds are perpendicular to the OCO plane and the electron lone pair orbital and $\sigma^{*}(\mathrm{CO})$ orbital strongly overlap. ${ }^{26}$ On the other hand, interactions between the OH and OMe bond dipoles
determine the energy difference between $\mathbf{5 a}$ (nearly antiparallel arrangement of bond dipoles, attraction) and 5b (parallel arrangement of bond dipoles, repulsion). The actual values of the dihedral angles OCOR $\left(\mathrm{R}=\mathrm{Me}\right.$ and $\mathrm{R}=\mathrm{H} ; \mathbf{5 a}, 69^{\circ}$ and $65^{\circ}$; 5b, $70^{\circ}$ and $-85^{\circ}$; Scheme 1a) are determined by both bond dipole interactions, anomeric delocalization, and secondorder hyperconjugation (bond staggering). ${ }^{27,28}$
Conformations 5a and 5b are connected by a flip-flop rotation (Scheme 1b) where, contrary to symmetric geminal double rotors such as $\mathrm{X}(\mathrm{OY})_{2}\left(\mathrm{X}=\mathrm{CH}_{2}, \mathrm{O}, \mathrm{S}\right.$, etc.; $\mathrm{Y}=\mathrm{CH}_{3}, \mathrm{OH}, \mathrm{SH}$, etc..$^{27,28}$ ), the flip-flop potential is characterized by two rather


5a $\quad \tau_{1}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{C})=68.9$
$\tau_{2}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{H})=64.9$

d $\quad \tau_{1}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{C})=4.2$
$\tau_{2}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{H})=80.4$


5b $\tau_{1}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{C})=69.9$
$\tau_{2}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{H})=-85.1$


5c $\tau 1(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{C})=85.3$
$\tau_{2}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{H})=-16.9$


Se $\begin{aligned} & \tau_{1}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{C})=180.0 \\ & \tau_{2}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{H})=56.4\end{aligned}$

$5 f \quad \tau_{1}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{C})=180$
$\tau_{2}(\mathrm{O}-\mathrm{C}-\mathrm{O}-\mathrm{H})=180$

Figure 2. Minimum conformations and rotational transition states of methoxymethanol (5) calculated at the B3LYP/6-311++G(d,p) level of theory. Distances in $\AA$, angles in deg. Dihedral angles $\tau_{1}$ and $\tau_{2}$ correspond to the convention used in Scheme 1.
than one rotational barrier. Rotation of the OH group (TS-5c; $\tau_{1}=85^{\circ}, \tau_{2}=-17^{\circ}$, Scheme 1b and Scheme 1a) is energetically less demanding (rotational barrier: $3.8 \mathrm{kcal} / \mathrm{mol}$ ) than rotation of the OMe group (TS-5d rotational barrier: $7.0 \mathrm{kcal} /$ $\mathrm{mol} ; \tau_{1}=4^{\circ}, \tau_{2}=80^{\circ}$; Scheme 1a). As indicated by the two barriers, the anomeric effect of the OMe group is clearly stronger than that of an OH group (reflected also by the two $-\mathrm{H}_{2} \mathrm{C}-$ OR bond lengths of $\mathbf{5 a}: \quad \mathrm{C}-\mathrm{OH}=1.412, \mathrm{C}-\mathrm{OMe}=1.402 \AA$; see Figure 2). For comparison, the barriers of rotation in methanol and $\mathbf{1}$ are 1.1 and $2.7 \mathrm{kcal} / \mathrm{mol}$, respectively. ${ }^{29}$

The flip-flop rotation of a geminal double rotor has been compared with the pseudorotation process of a five-membered ring. ${ }^{27,28}$ Accordingly, 5a corresponds to the twist form of a five-membered ring, while $\mathbf{5 b}$ is related to an envelope form. An "inversion" of $\mathbf{5 a}$ or $\mathbf{5 b}$ is an unlikely conformational process since it would involve simultaneous inward rotation of the two OR groups, thus leading to the sterically strongly hindered (sp,sp ) form, which occupies the global maximum of the conformational energy surface of $\mathbf{5}$. Another local minimum of the CES of $\mathbf{5}$ is associated with the "acyclic" form $\mathbf{5 e}(\mathrm{ap},+\mathrm{sc})\left(\tau_{1}\right.$ $=180^{\circ}, \tau_{2}=56^{\circ}$, Scheme 1 b and Scheme 1a), while the corresponding ( $+\mathrm{sc}, \mathrm{ap}$ ) form is just a transient point on the CES. Form $\mathbf{5 e}(2.2 \mathrm{kcal} / \mathrm{mol}$ relative to $\mathbf{5 a})$ benefits from an anomeric interaction of the OH group with the neighboring $\mathrm{C}-\mathrm{OMe}$ bond and from the anti arrangement with optimal bond staggering of the OMe group. The planar form TS-5f (ap,ap; Scheme 1a) is the transition state for the outward directed rotation of the OH group in $\mathbf{5 e}$ (barrier: $3.9 \mathrm{kcal} / \mathrm{mol}$ ).

Mechanism of the Formation of 5. The abstraction of an H atom from 1 by $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ is slightly exothermic. Utilizing the G2 method, ${ }^{30}$ we calculate reaction enthalpies $\Delta \Delta H_{\mathrm{f}}{ }^{0}(0)$ and $\Delta \Delta H_{\mathrm{f}}{ }^{0}-$ (298) for reaction 1 to be -5.4 and $-4.7 \mathrm{kcal} / \mathrm{mol}$, respectively. When these values are combined with the known heats of

$$
\begin{equation*}
\mathrm{CH}_{3} \mathrm{OCH}_{3}\left({ }^{1} \mathrm{~A}_{1}\right)+\mathrm{O}\left({ }^{3} \mathrm{P}\right) \rightarrow \mathrm{CH}_{3} \mathrm{OCH}_{2}\left({ }^{2} \mathrm{~A}\right)+\mathrm{OH}\left({ }^{2} \Pi\right) \tag{1}
\end{equation*}
$$

$$
\begin{equation*}
\mathrm{CH}_{3} \mathrm{OCH}_{2}\left({ }^{2} \mathrm{~A}\right)+\mathrm{OH}\left({ }^{2} \Pi\right) \rightarrow \mathrm{CH}_{3} \mathrm{OCH}_{2} \mathrm{OH}\left({ }^{1} \mathrm{~A}\right) \tag{2}
\end{equation*}
$$

formation for $\mathbf{1}, \mathrm{O}\left({ }^{3} \mathrm{P}\right)$, and the OH radical $(0 \mathrm{~K},-39.7,59.0$, $\left.9.2 ; 298 \mathrm{~K},-44.0,59.6,9.3 \mathrm{kcal} / \mathrm{mol}^{31}\right)$, the $\Delta H_{\mathrm{f}} 0(0)$ and $\Delta H_{\mathrm{f}} 0-$ (298) values of radical $\mathrm{CH}_{3} \mathrm{OCH}_{2}(6)$ are calculated to 4.7 and
$1.6 \mathrm{kcal} / \mathrm{mol}$, respectively. Various estimates of the heat of formation ranging from -6.9 to $-1.3 \pm 2 \mathrm{kcal} / \mathrm{mol}$ have been published. ${ }^{32}$ Good and Francisco ${ }^{33}$ have pointed out that these estimates of $\Delta H_{\mathrm{f}}{ }^{0}(298)$ are probably erroneous due to assumptions made in the experimental work. Their calculated $\Delta H_{\mathrm{f}}{ }^{0}(0)$ and $\Delta H_{\mathrm{f}}{ }^{0}(298)$ values of 4.2 and $0.9 \mathrm{kcal} / \mathrm{mol}$, respectively, are in reasonable agreement with our values and suggest that theory is more reliable in this case. ${ }^{33}$ The reaction enthalpies for H abstraction by $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ are $-50.8(0 \mathrm{~K})$ and $-50.1 \mathrm{kcal} / \mathrm{mol}(298$ K), respectively. The latter reaction should proceed without a barrier as is suggested by DFT and MP2 calculations as well as experimental observations. ${ }^{10}$ A collinear approach of $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ along one of the CH bond axes of $\mathbf{1}$ does not lead to a stationary point but leads via a bent $\mathrm{C} \cdots \mathrm{H} \cdots \mathrm{O}$ geometry to 5 as soon as linearity constraints are relieved in the geometry optimization. There is a possibility of probing the reactivity of the various CH bonds of $\mathbf{1}$ with the help of dimethyl ether oxide (modeling $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ attack) since H abstraction leads to a finite barrier in this case. It turns out that the out-of-plane H atoms can be more easily attacked than the in-plane H atoms (see below).

A computational investigation of H abstraction from 1 by O atoms faces both technical and methodological problems. The reaction mechanism can be complicated by the formation of complexes between $\mathbf{1}$ and $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ or $\mathbf{6}$ and OH . Little is known about the stability of complexes involving radicals; ${ }^{34}$ however, one can assume that the stability of these complexes is determined by the same forces as in the case of normal van der Waals complexes between closed-shell systems. On the other hand, the transition state energy for H abstraction by a radical will be low and, therefore, the potential energy surface (PES) along the path of the H abstraction will be rather flat. This means that the existence of a van der Waals complex as a precursor to the actual reaction complex will depend strongly on secondary effects such as zero-point energies (ZPE) and entropies rather than electronic effects. Therefore, the mechanistic relevance of complex formation in the H abstraction reaction is decided on the enthalpy $(\Delta H)$ or free enthalpy $(\Delta G)$ hypersurface (PHS or PGS) rather than the PES. There are examples in the literature which show that changes in ZPEs, temperature effects, and entropies rather than electronic changes along the reaction path can determine a reaction mechanism. ${ }^{35}$
1

7a



Figure 3. Geometries of dimethyl ether (1), $\mathrm{CH}_{3} \mathrm{OCH}_{2} \bullet$ radical (6), and stationary points 7 of the collision complexes between 1 and $\mathrm{O}(3 \mathrm{P})$ calculated at the UMPW1PW91/6-311 $++G(d, p)$ (normal print) and UMP2/6-311 $++G(d, p)$ level of theory (italics). For $\mathbf{1}$, the most likely directions of attack by $\mathrm{O}(3 \mathrm{P})$ are indicated by arrows.

DFT has been found to seriously fail with regard to the description of radical reactions such as $\mathrm{H}+\mathrm{H}_{2} \rightarrow \mathrm{H}_{2}+\mathrm{H},{ }^{36}$ $\mathrm{OH}+\mathrm{H}_{2} \rightarrow \mathrm{HOH}+\mathrm{H},{ }^{37} \mathrm{O}\left({ }^{3} \mathrm{P}\right)+\mathrm{CH}_{4} \rightarrow \mathrm{CH}_{3}+\mathrm{OH},{ }^{38}$ etc. This is mainly due to the self-interaction error ${ }^{39}$ caused by not using exact exchange in DFT as calculations with appropriate self-interaction correction (SIC) have shown. ${ }^{36,40}$ The SIC problem is partially avoided in hybrid DFT methods, which use exact exchange (at least partially). Hence, hybrid DFT methods provide a reasonable description of the energetics of reactions such as $\mathrm{OH}+\mathrm{H} \rightarrow \mathrm{HOH}+\mathrm{H} .{ }^{37}$ Also, B3LYP ${ }^{41}$ or the recently developed MPW1PW91 one-parameter hybrid functional ${ }^{42}$ are able to describe H -bonded complexes or in general van der Waals complexes better than standard DFT methods, provided a TZP basis with diffuse basis functions is used (see Experimental Section and refs 42 and 43). However, the rule of thumb is that the DFT description of a van der Waals complex is more problematic the more the complex stability is dominated by dispersion rather than electrostatic interactions. ${ }^{44}$ Taking this into consideration, we used unrestricted B3LYP and MPW1PW91 to explore the mechanism of reactions 1 and 2 and, then, we analyzed results with the help of UMP2 and UCCSD(T) calculations for the same compounds involved.

If one considers that exchange repulsion, polarization, and dispersion forces should dominate the interactions between $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ and $\mathbf{1}$, then complex formation should preferentially be expected in regions $3,2,4,5,6,1$, and 7 of $\mathbf{1}$ (see Figure 3), since in this order exchange repulsion will increase while stabilizing dispersion and polarization forces will decrease. Actually, explorative calculations with a DZP basis in these regions surrounding 1 led to the discovery of three complexes, three first-order transition states, and two second-order transition states. Utilizing these stationary points for refined descriptions with the $6-311++G(d, p)$ basis at UB3LYP and UMPW1PW91 structures $7 \mathbf{a}$ (complex between $\mathbf{1}$ and $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$, 7b (transition state), 7c (complex between 6 and $\bullet \mathrm{OH}$ ), and other structures such as $7 \mathbf{d}$ (complex between 1 and $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$, see Figure 3 and Table 7) were found. The corresponding energies, enthalpies, and free enthalpies are given in Table 6.

Clearly, $\mathbf{7 a}, \mathbf{7 b}$, and $\mathbf{7 c}$ are located at three consecutive stationary points along the path of reaction 1 involving one of the out-of-plane H atoms. A similar approach of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ in the direction of the bond axis of CH (in-plane) involves a secondorder TS, which is about $8 \mathrm{kcal} / \mathrm{mol}$ higher in energy than TS 7b. An adiabatic analysis ${ }^{45,46}$ of the B3LYP/6-311++G(d,p) vibrational modes of $\mathbf{1}$ leads to adiabatic CH stretching frequencies of $3102 \mathrm{~cm}^{-1}$ (in plane) and $2980 \mathrm{~cm}^{-1}$ (out-ofplane) $($ force constants, $k$ (in-plane) $=5.271, k$ (out-of-plane $)$ $=4.864 \mathrm{mdyn} / \AA ̊$; exptl stretching frequencies, 2981, 2874 $\mathrm{cm}^{-1}$ ), which according to the relationships given by Larsson and Cremer ${ }^{47}$ corresponds to CH dissociation energies $D_{\mathrm{e}}=$ 112.4 and $102.3 \mathrm{kcal} / \mathrm{mol}\left(D_{0}=103.8\right.$ and $\left.93.7 \mathrm{kcal} / \mathrm{mol}\right)$, respectively. This indicates that abstraction of an out-of-plane H atom should be $10 \mathrm{kcal} / \mathrm{mol}$ easier than that of an in-plane H atom of $\mathbf{1}$ in line with the DFT results obtained for the reaction of 1 with $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ or $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$. Also, abstraction of an out-of-plane H atom leads to the more stable $\mathrm{CH}_{3} \mathrm{OCH}_{2} \bullet$ form, as can be seen by comparing the geometries of $\mathbf{7 a}, \mathbf{7 b}$, and $\mathbf{7 c}$ with that of radical 6 (see Figure 3 and Table 7). While these considerations apply to a fictitious molecule $\mathbf{1}$ at rest, one has to consider that in the real situation the methyl groups of $\mathbf{1}$ are rapidly rotating. Hence, an attack of an O atom at the in-plane H will lead to CH bond rupture and H transfer if the O atom can follow the movement of the H atom in an out-of-plane position where a rupture of the CH bonds becomes easier.

The PES is rather flat for an approach of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ toward one of the out-of-plane H atoms. At UB3LYP/6-311++(d,p), complex $7 \mathbf{a}$ is $1.7 \mathrm{kcal} / \mathrm{mol}$ (after BSSE corrections; $2.0 \mathrm{kcal} /$ mol after additional ZPE corrections) below the separated reactants, but its stability decreases to $0.1 \mathrm{kcal} / \mathrm{mol}$ for the more reliable MPW1PW91 hybrid functional. At the UMP2/6$311++(\mathrm{d}, \mathrm{p})$ level, 7a no longer exists, which clearly shows that UB3LYP exaggerates stabilizing O,H interactions at short O10, H8 distances (Figure 3 and Table 7) and leads to spurious van der Waals complexes such as 7d. This tendency decreases when using the MPW1PW91 functional, but nevertheless it is also a

TABLE 6: Energetics of Reaction 1 Calculated at Various Levels of Theory ${ }^{a}$

| compd | $\Delta E$ | $\Delta E(\mathrm{BSSE})$ | ZPE | $E_{0}$ | $\Delta H(298)$ | $\Delta G(298)$ | $\mu$ | $\omega$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| I. B3LYP/6-31G(d,p) |  |  |  |  |  |  |  |  |
| $1+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ | -230.09359 | 0 | 50.1 | 0 | -230.00608 | -230.05335 | 1.27 | 221 |
| 7a | -2.7 | -1.1 | 50.3 | -1.0 | -1.2 | 4.4 | 0.95 | 73 |
| 7b | -1.4 | -0.1 | 47.2 | -3.1 | -3.6 | 3.5 | 2.17 | 888i |
| 7c | -9.2 | -7.7 | 48.9 | -8.9 | -9.0 | -3.4 | 2.18 | 101 |
| $6+\mathrm{OH}$ | -2.4 | -2.4 | 46.7 | -5.8 | -5.2 | -7.8 | 1.35 | 171 |
| barrier | 1.2 | 1.0 | - |  |  | -3.5 |  |  |
| stability of 7c | -6.8 | -5.2 | - | -3.1 | -3.8 |  |  |  |
| II. B3LYP/6-311++G(d,p) |  |  |  |  |  |  |  |  |
| $1+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ | -230.16699 | 0 | 49.7 | 0 | -230.08010 | -230.12750 | 1.44 | 205 |
| 7 a | -2.0 | -1.7 | 49.3 | -2.0 | -2.2 | 3.1 | 1.65 | 47 |
| 7b | -1.7 | -1.4 | 47.2 | -3.8 | -4.6 | 2.3 | 2.67 | 442i |
| 7c | -10.8 | $-10.5$ | 48.1 | -12.2 | -12.0 | -7.5 | 2.96 | 37 |
| $6+\mathrm{OH}^{\bullet}$ | -6.2 | -6.3 | 46.3 | -9.7 | -9.3 | -11.7 | 1.50 | 159 |
| barrier | 0.3 | 0.3 |  |  |  |  |  |  |
| stability of 7c | -4.3 | -4.2 |  | $-2.5$ | -2.6 |  |  |  |
| III. MPW1PW91/6-311++(d,p) |  |  |  |  |  |  |  |  |
| $1+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ | -230.0990 | 0 | 50.2 | 0 | -230.01142 | -230.05878 | 1.42 | 207 |
| 7 a | -0.4 | -0.1 | 45.0 | -0.3 | -0.3 | 3.9 | 1.22 | 17 |
| 7b | 0.4 | 0.7 | 47.5 | -1.9 | -2.7 | 4.3 | 2.56 | 547i |
| 7c | -9.4 | -9.1 | 48.5 | -10.7 | -10.5 | -6.1 | 2.90 | 37 |
| $6+\mathrm{OH}^{+}$ | -5.0 | -5.1 | 46.8 | -8.5 | -7.8 | -10.1 | 1.45 | 163 |
| barrier | 0.7 | 0.8 |  |  |  | -4.3 |  |  |
| stability of 7c | -4.4 | -4.0 |  | $-2.3$ | $-2.7$ | -4.0 |  |  |
| IV. MP2(frozen)/6-311++G(d,p) |  |  |  |  |  |  |  |  |
| $\mathbf{1}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ | -229.54715 | 0 | 50.8 | 0 | -229.45703 | -229.50434 | 1.65 | 198 |
| 7b (barrier) | 14.9 | 15.2 | 47.6 | 12.1 | 10.6 | 17.2 | 1.51 | 2077i |
| 7c | -2.1 | -1.7 | 49.5 | -3.0 | -3.7 | 0.0 | 2.77 | 41 |
| $6+\mathrm{OH}$ | 0.2 | $-0.3$ | 47.6 | -3.6 | -2.8 | -5.3 | 1.56 | 169 |
| barrier | 14.2 | 15.2 |  | 12.1 | 10.6 | 17.2 |  |  |
| stability of 7c | $-2.3$ | -1.3 |  |  | -0.9 |  |  |  |
| V. $\operatorname{CCSD}(\mathrm{T}$, frozen)/6-311++G(d,p)//MP2(frozen)/6-311++G(d,p) |  |  |  |  |  |  |  |  |
| $1+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ | -229.61362 | 0 | 50.8 | 0 | $-229.52350$ | -229.57081 |  |  |
| 7b (barrier) | 10.2 | 10.6 | 47.6 | 7.4 | 5.9 | 12.5 |  |  |
| 7c | -2.8 | -2.4 | 49.5 | -3.7 | -4.4 | -0.7 |  |  |
| $6+\mathrm{OH}^{\bullet}$ | 0.8 | 0.3 | 47.6 | -2.9 | -2.1 | -4.6 |  |  |
| barrier (corrected) ${ }^{b}$ | 7.6-t | $8.0-\mathrm{t}$ |  | 4.8-t | 2.3 | 8.9 |  | $E_{\text {a }}=3.5$ |
| stability of 7c | -3.6 | -2.7 |  | -0.8 | -2.3 |  |  |  |

${ }^{a}$ Absolute energies of reference system $1+\mathrm{O}(3 \mathrm{P})$ in hartree, relative energies values in $\mathrm{kcal} / \mathrm{mol} . \Delta(\mathrm{BSSE})$, energy differences after basis set superposition corrections; ZPE, zero-pint corrections; $E_{0}$, energy differences at 0 K after ZPE corrections; $\Delta H(298)$ and $\Delta G(298)$, enthalpy and free enthalpy differences at $298 \mathrm{~K} ; \mu$, dipole moments in Debye; $\omega$, lowest vibrational frequency in $\mathrm{cm}^{-1}$. For UCCSD(T) results, vibrational and temperature corrections were taken from UMP2 calculations. ${ }^{b}$ Corrections comprise the G2 correction of $-2.6 \mathrm{kcal} / \mathrm{mol}$ and the Wigner tunneling correction of $-0.96 \mathrm{kcal} / \mathrm{mol}$ at 298 K . See text. Tunneling corrections at 8 K can no longer calculated by the Wigner formula, and therefore, the correction due to tunneling is indicated by $\mathrm{t} . E_{\mathrm{a}}$ is the calculated Arrhenius activation energy at 298 K .
basic deficiency of the new functional, which was actually designed to improve the description of van der Waals complexes. ${ }^{42}$

As a direct consequence of exaggerating bonding interactions between O 10 and H 8 , the reaction barrier (TS-7b) is underestimated by about $6 \mathrm{kcal} / \mathrm{mol}$ (see below), where it does not matter whether UB3LYP ( $0.3 \mathrm{kcal} / \mathrm{mol}$ ) or UMPW1PW91 ( 0.8 $\mathrm{kcal} / \mathrm{mol}$ ), a large or a small basis set, is used (Table 6). UMP2 exaggerates at the same time the barrier by about $8 \mathrm{kcal} / \mathrm{mol}$, suggesting a value of $15 \mathrm{kcal} / \mathrm{mol}$. Despite the large difference in calculated energies, both UDFT and UMP2 suggest similar geometries for TS-7b. The O10, H8 distance varies between $1.37 \AA$ (UDFT) and $1.30 \AA$ (UMP2), while the C3, H8 distance is predicted to be $1.2 \AA$ by all three methods (Figure 3). Calculated geometries are qualitatively in line with what one can expect for the geometry of TS-7b.

While it is presently not possible to improve the description of a TS such as 7b at the DFT level, more reliable calculations can be employed at the ab initio level. $\operatorname{UCCSD}(\mathrm{T}) / 6-311++\mathrm{G}-$ (d,p) predicts a barrier of $10.6 \mathrm{kcal} / \mathrm{mol}$ at the UMP2 geometry. The corresponding activation enthalpy at 298 K is $5.9 \mathrm{kcal} /$ mol , and the free activation energy is $12.5 \mathrm{kcal} / \mathrm{mol}$. These
values have to be corrected since tunneling should play an important role for reaction 1. At 298 K , a Wigner correction ${ }^{48}$ of $0.96 \mathrm{kcal} / \mathrm{mol}$ is obtained, which leads to a tunneling corrected activation enthalpy $\Delta H^{\#}$ of $4.9 \mathrm{kcal} / \mathrm{mol}$. From G2 calculations, we can estimate that an increase of the basis set and higher order correlation effects will lead to an additional lowering of the barrier by $2.6 \mathrm{kcal} / \mathrm{mol}$, which is similar to the error in the calculated $\operatorname{UCCSD}(\mathrm{T})$ heat of reaction at $298 \mathrm{~K}(-2.1 \mathrm{kcal} /$ mol compared to the G 2 value of $-4.7 \mathrm{kcal} / \mathrm{mol}$; see Table 6). This is in line with results of previous investigations on the system $\mathrm{OH}+\mathrm{H}_{2} \rightarrow \mathrm{HOH}+\mathrm{H}$, which showed that geometry optimization of a TS such as 7b at $\operatorname{UCCSD}(\mathrm{T})$ with a QZ3P( 2 f 1 g ) basis set can lower the barrier by about $2-3 \mathrm{kcal} / \mathrm{mol}^{49}$ Hence, the final $\Delta H^{\#}(298)$ obtained in this way is $2.3 \mathrm{kcal} /$ mol.

LeFevre and co-workers obtained an Arrhenius activation energy of $E_{\mathrm{a}}=2.85 \pm 0.5 \mathrm{kcal} / \mathrm{mol}$ by combining massspectrometric and ESR results. ${ }^{8}$ Since $E_{\mathrm{a}}=\Delta H^{\#}+2 R T$ for reaction 1, the experimental value compares well with the $\operatorname{UCCSD}(\mathrm{T})$ estimate of $E_{\mathrm{a}}=3.5 \mathrm{kcal} / \mathrm{mol}$. As for the reaction barrier at 8 K (temperature of reaction in the matrix), we obtain a value of $7.4 \mathrm{kcal} / \mathrm{mol}$, which has to be corrected by $2.6-4.8$

TABLE 7: Calculated Geometries of Stationary Points Encountered in Reaction 1 (Compare with Figure 3) ${ }^{a}$


[^1]$\mathrm{kcal} / \mathrm{mol}$. Assuming that the calculated temperature corrections are reasonable, the corresponding experimental value should be $4.2 \mathrm{kcal} / \mathrm{mol}$. At this temperature, tunneling is large (but can no longer be estimated by the Wigner formula ${ }^{48 b}$ ) so that the barrier becomes effectively zero.

In the exit channel of reaction 1, both UDFT, UMP2, and $\operatorname{UCCSD}(\mathrm{T})$ predict the existence of a radical-radical complex 7c formed by 6 and OH in a triplet state. The predicted stability of complex $\mathbf{7 c}$ ranges from $1.3 \mathrm{kcal} / \mathrm{mol}$ (UMP2) to maximally $5.2 \mathrm{kcal} / \mathrm{mol}$ (Table 6), where again DFT seems to exaggerate stabilizing interactions between the radical partners. Inclusion of ZPE corrections reduces the complex stability to 0.8 $(\operatorname{UCCSD}(\mathrm{T}))$ and $2.5 \mathrm{kcal} / \mathrm{mol}(\mathrm{UB} 3 \mathrm{LYP} / 6-311++\mathrm{G}(\mathrm{d}, \mathrm{p}))$, while the calculated enthalpy differences at 298 K suggest a stable complex also on the PHS (stability: $0.9-3.5 \mathrm{kcal} / \mathrm{mol}$; Table 6) irrespective of the method and basis set used. Entropy destroys the complex as is suggested by calculated $\Delta G(298)$ values (Table 6).

The various geometries calculated for $7 \mathbf{c}$ show the same relatively small variation in dependence of the method and basis set used as was found for TS-7b (Figure 3 and Table 7). At UMP2, the OH bond length of the donor molecule OH is lengthened by $0.008 \AA$ similarly as one has found for the $\mathrm{H}_{2} \mathrm{O}$ dimer $(0.006 \AA) .{ }^{50}$ Considering that a CH bond $(1.09 \AA)$ is about
$0.14 \AA$ longer than a OH bond $(0.95 \AA)$, the OC distance should be about $3.12 \AA$ if one uses the experimental OO distance ( 2.98 $\AA$ ) in the water dimer. ${ }^{50}$ The calculated OC distance is $3.15 \AA$ (UMP2) but somewhat smaller when UDFT is applied (3.00$3.08 \AA$, Table 7), which reflects the exaggeration of stabilizing interactions at DFT. The stability of H -bond complex 7c $(\operatorname{UCCSD}(\mathrm{T}): 2.7 \mathrm{kcal} / \mathrm{mol}$, Table 6) is about half of that of the water dimer $(5 \mathrm{kcal} / \mathrm{mol}) .{ }^{50}$ Clearly, theory suggests the existence of an H bonded triplet complex between OH and $\mathbf{6}$ where the donor is a radical and the acceptor is a C atom with a single electron. Such complexes should also exist in other reactions involving $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$; however, it will be difficult to detect these complexes since the excess energy such a complex will possess (about $5 \mathrm{kcal} / \mathrm{mol}$ if one considers a barrier of $2.3 \mathrm{kcal} / \mathrm{mol}$; see above) should be higher than the complex stability.

Despite the somewhat different descriptions provided by $\operatorname{UDFT}, \operatorname{UCCSD}(\mathrm{T})$, and UMP2, all theories agree with regard to predictions relevant for the experiments carried out at 8 K in this work. The energetically most favorable attack of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ is at one of the out-of-plane H atoms of $\mathbf{1}$, and there is no energy barrier since H tunneling will reduce a barrier of $4.8 \mathrm{kcal} / \mathrm{mol}$ effectively to zero.

Recombination of radicals 6 and OH (reaction 2) should occur within a cavity of the matrix at 10 K without a barrier. Formation

## SCHEME 2. Recombination of $\mathrm{CH}_{3} \mathbf{O C H}_{2}$ and $\mathbf{O H}$ Radical Assuming the Geometry of Complex $7 \mathbf{c}^{a}$


${ }^{a}$ Rotation a is sterically less hindered and leads to $\mathbf{5 A}$, rotation $b$ to $\mathbf{5 b}$.
of $\mathbf{5}$ is strongly exothermic $\left(\Delta \Delta H_{\mathrm{f}}{ }^{0}(0)\right.$ and $\Delta \Delta H_{\mathrm{f}}{ }^{0}(298)$ : -98.5, $-100.1 \mathrm{kcal} / \mathrm{mol}, \mathrm{G} 2$ calculations). The heats of formation $\Delta H_{\mathrm{f}}{ }^{0}-$ (O) and $\Delta H_{\mathrm{f}}{ }^{0}(298)$ of $\mathbf{5}$ can be calculated from reactions 1 and $2\left(\Delta \Delta H_{\mathrm{f}}{ }^{0}(0)\right.$ and $\Delta \Delta H_{\mathrm{f}}{ }^{0}(298),-103.9,-104.8 \mathrm{kcal} / \mathrm{mol}$; values for $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ insertion into $\left.\mathbf{1}-149.3,-150.2 \mathrm{kcal} / \mathrm{mol}\right)$ to be -84.6 and $-89.2 \mathrm{kcal} / \mathrm{mol}$, respectively, which agrees well with an $\Delta H_{\mathrm{f}}{ }^{0}(298,5)$ estimate of -89.4 based on thermochemical group increments. ${ }^{51}$

It is interesting to note that radical $\mathbf{6}$ is no longer a pure $\pi$-radical. Pyramidalization at the radical center is given by the difference $360-\Sigma\left(\mathrm{C}\right.$ bond angles) $=8.5^{\circ}$ (for the geometry, see Figure 3 and Table 7). Clearly, this is a result of destabilizing three-electron interactions involving a single $\pi$ electron at C and the two electrons occupying a $\mathrm{p} \pi$ lone pair orbital at O . Theory shows that this can be avoided by (a) reducing the overlap between the orbitals involved and (b) increasing the difference between the corresponding orbital energies. Hence, the $\sigma$ nature of radical $\mathbf{6}$ is in line with observations made for $\bullet \mathrm{CH}_{2} \mathrm{X}(\mathrm{X}=\mathrm{OH}, \mathrm{F}, \mathrm{Cl})$ radicals. ${ }^{52}$

In TS-7b as well as complex 7c, the equilibrium geometry of radical $\mathbf{6}$ is largely developed. If the OH group rotates as shown in Scheme 2 where the rotation may be initiated by a large amplitude OHC bending vibration, then, depending on the direction of rotation, conformations $\mathbf{5 a}$ or $\mathbf{5 b}$ will be directly formed. Steric interactions should be smaller for rotation (a) in Scheme 2 (thus leading to more 5a); however, this depends on the actual distance between $\mathbf{6}$ and OH , for which rotation will become possible.

## Conclusion

The oxidation of dimethyl ether $\mathbf{1}$ with atomic oxygen generated by photolysis of ozone or $\mathrm{N}_{2} \mathrm{O}$ has been examined in low-temperature matrixes. The major reaction products are the two most stable conformers of methoxymethanol, $\mathbf{5 a}$ and $\mathbf{5 b}$, while the third conformation $\mathbf{5 e}$ was not found among the reaction products. The weakly bound ( 2.8 kcal ) charge transfer complex of dimethyl ether with atomic oxygen $\mathrm{O}\left({ }^{3} \mathrm{P}\right), \mathbf{3}$, was also not observed. Experience collected in this work suggests that such a complex does not exist and is just an artifact of the DFT calculation used in ref 9 .

Three mechanisms have to be considered to rationalize the formation of 5 .
(i) The most reasonable mechanism is the hydrogen abstraction from 1 to give $\mathbf{6}$ and OH radicals, which by in-cage recombination produce 5. According to a calculated heat of formation of $4.7 \mathrm{kcal} / \mathrm{mol}$ (at 0 K ), radical $\mathbf{6}$ is quite stable, however, due to a rapid in-cage reaction not observed in the
experiments. This mechanism requires that the small OH radical does not escape the matrix cage in significant amounts.
(ii) The second mechanism is the direct insertion of O (1D) into a CH bond of ether $\mathbf{1}$ without formation of intermediate radicals.
(iii) The third mechanism $\mathbf{1}$ is the rearrangement of dimethyl ether O-oxide $\mathbf{2}$ formed as a short-lived intermediate. The barrier for the fragmentation of $\mathbf{2}$ into $\mathbf{6}$ and OH radicals is predicted to be $25.9 \mathrm{kcal} / \mathrm{mol}$, and these radicals could recombine to 5 as described above. However, with an activation barrier of only $16 \mathrm{kcal} / \mathrm{mol}$, the fragmentation into $\mathrm{CH}_{3}$ and $\mathrm{CH}_{3} \mathrm{OO}$ represents the lowest exit channel for $2 .{ }^{9}$ Neither of these radicals nor their recombination product dimethylperoxide 4 could be observed in the IR spectrum. Thus, the route via dimethyl ether O-oxide 2 can clearly be excluded as a possible mechanism for the oxidation of ether 1.

Possibilities (i) or (ii) are both likely as reaction mechanism. H abstraction by $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ occurs preferentially at the out-of-plane CH bonds of $\mathbf{1}$ and in this case the activation energy of reaction 1 is effectively zero. The calculated Arrhenius energy of 3.5 $\mathrm{kcal} / \mathrm{mol}$ at 298 is in agreement with an experimental $E_{\mathrm{a}}$ of 2.85 $\pm 0.5 \mathrm{kcal} / \mathrm{mol}$.

Stabilization of radical $\mathbf{6}$ by three-electron interactions, which are already active in TS-7b lower the barrier by about $2 \mathrm{kcal} /$ mol as can be deduced from the activation enthalpies measured for the corresponding reaction in the case of alkanes (propane $\left.+\mathrm{O}\left({ }^{3} \mathrm{P}\right): 5 \mathrm{kcal} / \mathrm{mol}\right) .{ }^{53}$

As an unexpected result of the investigation of the reaction mechanism, we found a H -bonded radical-radical complex in a triplet state that involves C as an acceptor. To the best of our knowledge, such complexes have never been described before and, although their experimental verification may be difficult, they should play an important role in the elucidation of the reaction mechanism involving radicals.

## Experimental Section

Matrix Spectroscopy. Matrix isolation experiments were performed by standard techniques with an APD CSW-202 Displex closed cycle helium cryostat. Matrixes were produced by deposition of argon (Messer-Griesheim, 99.9999\%), mixtures of argon, dimethyl ether, and ozone, or mixtures of argon, dimethyl ether, and nitrous oxide on top of a CsI (IR) or sapphire (UV-vis) window with a rate of approximately $0.15 \mathrm{mmol} /$ min. To obtain optically clear matrixes, the cold window was retained at $25-30 \mathrm{~K}$ (Ar/ether/ozone) or 17 K (Ar/ether/N2O) during deposition and afterward cooled to $7-9 \mathrm{~K}$.

Infrared spectra were recorded by using a Bruker IFS66 FTIR spectrometer with a standard resolution of $1 \mathrm{~cm}^{-1}$ in the range
of $400-4000 \mathrm{~cm}^{-1}$. Irradiations were carried out with use of a KrF Excimer Laser $(\lambda=248 \mathrm{~nm})$ or ArF Excimer Laser $(\lambda=$ 193 nm ) (LPX 105 SD Lambda-Physic) or with use of an Osram HBO $500 \mathrm{~W} / 2$ mercury high-pressure arc lamp in Oriel housings equipped with quartz optics. IR irradiation from the lamps was absorbed by a 10 cm path of water. Schott cut off filters were used ( $50 \%$ transmission at the wavelength specified) in combination with dichroic mirrors.

Computational Details. Investigations were carried out employing different levels of theory. Restricted and unrestricted Kohn-Sham theory was applied using two different hybrid functionals, namely Becke's three-parameter functional B3LYP ${ }^{41}$ and the modified Perdew-Wang functional described in ref 42, which is supposed to describe van der Waals complexes more accurately. Pople's $6-31 \mathrm{G}(\mathrm{d}, \mathrm{p})$ and $6-311++\mathrm{G}(\mathrm{d}, \mathrm{p})$ basis sets ${ }^{54}$ were employed to determine geometry, conformation, and vibrational spectra of $\mathbf{1 , 5}, \mathbf{6}$, and $\mathbf{7}$, where the VTZ basis with the diffuse functions was needed in particular for the description of molecular complexes formed in reactions 1 and 2 . For the investigation of the latter as well as the insertion of $\mathrm{O}\left({ }^{1} \mathrm{D}\right)$ into 1, B3LYP and MPW1PW91 results were checked by employing unrestricted MP2 and MP4 perturbation theory ${ }^{56}$ as well as unrestricted $\operatorname{CCSD}(\mathrm{T})$ theory where for MP4 and $\operatorname{CCSD}(\mathrm{T})$ only single-point calculations were carried out. For the radicals investigated, spin contamination was determined by calculating the expectation value of $S^{2}$, which in no case deviated from the ideal values by more than $10 \%$.

For all molecules and transition states considered, vibrational frequencies at optimized geometries were determined to characterize stationary points. Zero-point energy (ZPE) and thermal corrections were used to obtain reaction and activation enthalpies at 298 K . The determination of entropies led to free enthalpies at 298 K . Some of these calculations were carried out for nonstationary points. In these cases, the determination of vibrational corrections was based on generalized vibrational modes. ${ }^{57}$ Reaction energies and enthalpies of reactions 1 and 2 were determined with the help of G2 theory according to standard procedures described in the literature. ${ }^{30}$ Utilizing experimental heats of formation $\Delta H_{\mathrm{f}}{ }^{0}(298)$ for suitable reference compounds, ${ }^{31}$ heats of formation for radical 6 and closed-shell molecule 5 were determined.

For the van der Waals complexes 7, basis set superposition errors (BSSE) were corrected with the help of the counterpoise method. ${ }^{58}$ Calculations were carried out with COLOGNE96, ${ }^{59}$ ACES, ${ }^{60}$ GAUSSIAN94, ${ }^{61}$ and GAUSSIAN98. ${ }^{62}$

Materials. Mass spectra (EI, 70 eV ) were taken on a Varian MAT CH5.

Dimethyl ether 1 was obtained from Aldrich Company ( $99 \%$ purity) and used without further purification. Dimethyl ether180 was prepared by a procedure described by Blukis, Kasai, and Myers. ${ }^{63}$ By mass spectroscopy the ${ }^{18} \mathrm{O}$ content was determined to be $86 \%$. Dimethyl ether $-d_{6}$ was prepared by a procedure described by Kanazawa and Nukada, ${ }^{64}$ and the deuterium content was determined to be $99 \%$. Ozone and Ozone${ }^{18} \mathrm{O}$ were generated with an ozonizer (Demag), trapped at 77 K , and purified according to ref 65 . Ozone- ${ }^{18} \mathrm{O}$ was obtained by using ${ }^{18} \mathrm{O}_{2}$ (99.5\% isotopic purity) in the procedure described above. Nitrous oxid $\mathrm{N}_{2} \mathrm{O}$ was obtained from Air Liquide Company ( $99.9 \%$ purity).

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Supporting Information Available: Z-Matrixes of compounds $\mathbf{1 , 5} \mathbf{5}$, and $\mathbf{7}$ including various conformations and the corresponding transition states. This material is available free of charge via the Internet at http://pubs.acs.org.

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[^1]:    ${ }^{a}$ Basis sets are denoted in short form: $\mathrm{A}=6-31 \mathrm{G}(\mathrm{d}, \mathrm{p}) ; \mathrm{B}=6-311++\mathrm{G}(\mathrm{d}, \mathrm{p})$. All bond lengths in $\AA$, bond angles in deg. ${ }^{b}$ For 7a, the UMP2 optimization led to a second-order TS with two imaginary frequencies

